

## Exotic Structures of $\text{Si}_2\text{B}_2\text{H}_4$

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The potential energy surface of  $\text{Si}_2\text{B}_2\text{H}_4$ , explored *ab initio* at the MP2/6-31G\* and density functional theory, Becke3LYP/6-311+G\*\* levels, is rich in detail. Twelve minima were found out of the 39 stationary points located. 1,3-Disiladiboretene (**1a**), with a puckered four-membered ring, is predicted to be the global minimum. Its calculated inversion barrier (1.2 kcal/mol), however, is much lower than that of the analogous 1,3-diboretene (17.0 kcal/mol). The relative energies of two other  $\text{Si}_2\text{B}_2\text{H}_4$  isomers (**5** and **7**) with B–H–Si bridges are only 2.0 kcal/mol above **1a**. The relative energies of **9** and **15a**, both with planar pentacoordinate borons, are 7.3 and 27.1 kcal/mol. The acyclic (**18**) and linear (**22a,b**) isomers are noncompetitive energetically. Many other  $\text{Si}_2\text{B}_2\text{H}_4$  silaboranes, possessing similar structures and bonding as the corresponding carboranes, are calculated to be saddle points. In general, hydrogen bridging and cyclic  $\pi$ -delocalization stabilize  $\text{Si}_2\text{B}_2\text{H}_4$  isomers, while structures with multiply bonded silicon are unfavorable.

## Introduction

Theoretical calculations first revealed the cyclobutadiene dication ( $\text{C}_4\text{H}_4^{2+}$ )<sup>1</sup> to be nonplanar (Figure 1) despite being Hückel  $2\pi$  electron aromatic.<sup>2</sup> Subsequently, the excellent agreement between the computed NMR-IGLO chemical shifts<sup>3</sup> and the experimental values<sup>4</sup> provided conclusive evidence for the nonplanarity of the tetramethyl derivative  $\text{C}_4(\text{CH}_3)_4^{2+}$ . The neutral isoelectronic 1,3-diboretene ( $1,3\text{-C}_2\text{B}_2\text{H}_4$ ) and  $1,3\text{-C}_2\text{B}_2\text{H}_2(\text{NH}_2)_2$  also have nonplanar geometries and similar electronic structures.<sup>5</sup> Moreover, the theoretical predictions were verified subsequently by X-ray crystallography on  $1,3\text{-t-Bu}_2\text{C}_2(\text{BNMe}_2)_2$ <sup>6b</sup> and on other derivatives.<sup>6a,c–g</sup> However, Korkin *et al.*'s<sup>7a</sup> recent computations found  $\text{Si}_4\text{H}_4^{2+}$ , the homologue of  $\text{C}_4\text{H}_4^{2+}$ , to be nearly planar.<sup>7</sup> What is the situation with 1,3-disiladiboretene ( $1,3\text{-Si}_2\text{B}_2\text{H}_4$ )?

While the potential energy surfaces (PES) of  $\text{C}_4\text{H}_4^{2+}$  and  $\text{Si}_4\text{H}_4^{2+}$  have not been explored fully,<sup>1b,7</sup> *ab initio* MO studies revealed remarkable structural variations for  $\text{C}_2\text{B}_2\text{H}_4$  isomers.<sup>8</sup> The resulting fruitful interplay between theory and experiment guided the synthesis of many  $\text{C}_2\text{B}_2\text{R}_4$  isomers by Berndt and co-workers.<sup>9</sup> Such successes serve as additional stimuli for the present exploration of the PES of  $\text{Si}_2\text{B}_2\text{H}_4$ .

The experimental literature on silaboranes is rather scarce.<sup>10</sup> The recent synthesis and X-ray structure determination of icosahedral mono- and disilaboranes by Seyferth, Wesemann, and their co-workers<sup>10</sup> mark the beginning of the combined boron–silicon chemistry. Smaller silaboranes like  $\text{SiBH}_n$  ( $n = 1–5$ ),<sup>11</sup>  $\text{Si}_2\text{B}_2\text{H}_4$ ,<sup>12a</sup> and  $\text{Si}_2\text{B}_3\text{H}_3$ <sup>12b</sup> have been explored only computationally. These studies point to a structural correlation between boranes and silaboranes, i.e. the analogy between trivalent boron and divalent silicon moieties.<sup>13</sup> Thus, by equating a B–H bond pair to Si: ( $\sigma$  lone pair), the doubly bridged  $\text{Si}(\mu\text{-H})_2\text{Si}$  ( $\text{C}_{2v}$ )<sup>14</sup> and silacycloprenylidene ( $\text{C}_2\text{H}_2\text{-Si}$ )<sup>15a,c</sup> are related to the doubly bridged  $\text{B}_2\text{H}_4$  ( $\text{HB}(\mu\text{-H})_2\text{BH}$ ,  $\text{C}_{2v}$ ) and to borirene ( $\text{C}_2\text{H}_2\text{BH}$ ,  $\text{C}_{2v}$ ).<sup>16</sup>

Further *ab initio* studies on small silicon compounds<sup>17</sup> have

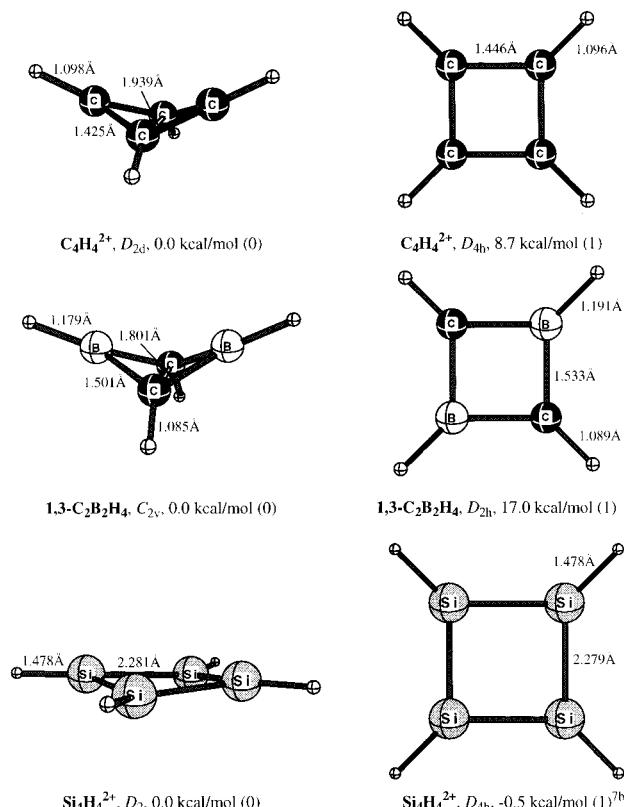


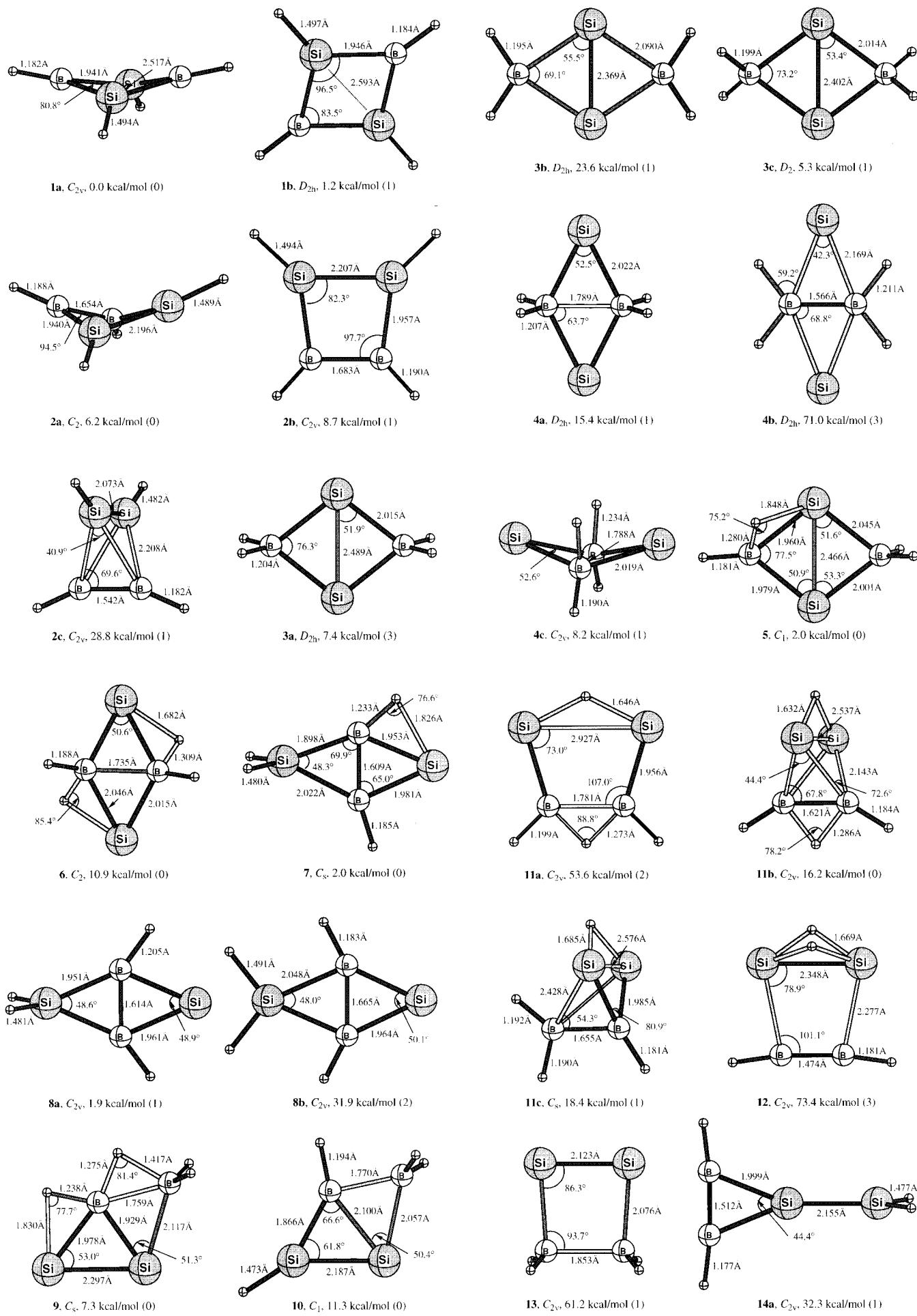
Figure 1. Becke3LYP/6-311+G\*\* optimized geometries of  $\text{C}_4\text{H}_4^{2+}$ ,  $1,3\text{-C}_2\text{B}_2\text{H}_4$ , and  $\text{Si}_4\text{H}_4^{2+}$ .

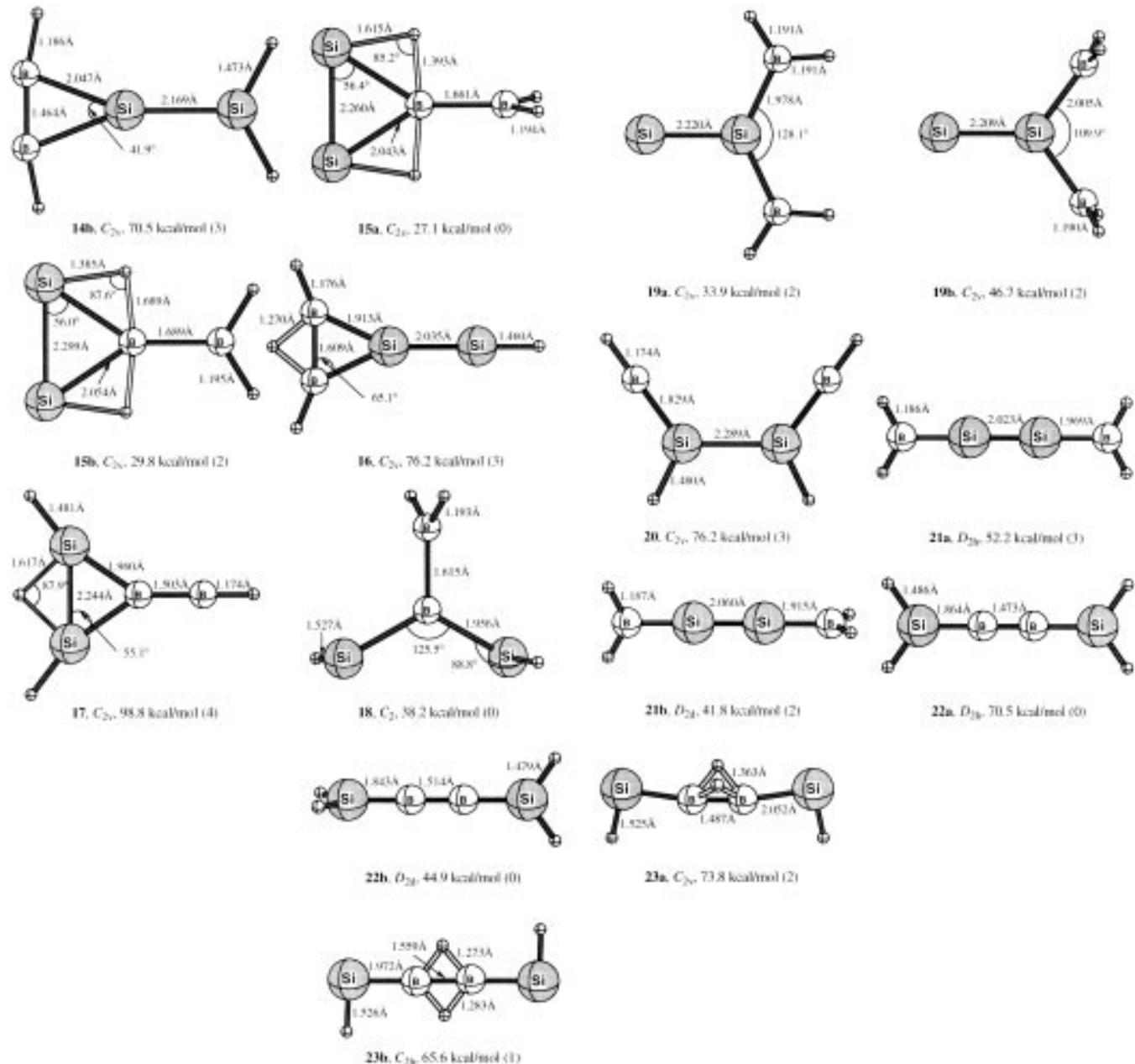
shown that the structural preferences can be quite different from those of the corresponding carbon compounds. For example,  $\text{Si}(\mu\text{-H})_2\text{Si}$  was predicted to be the global minimum on the PES of  $\text{Si}_2\text{H}_2$ <sup>17c</sup> rather than  $\text{HSi}\equiv\text{SiH}$ , and this was verified spectroscopically.<sup>14</sup> Similarly, silylsilylene ( $\text{HSi}-\text{SiH}_3$ ) and  $\text{HSi}(\mu\text{-H})_2\text{SiH}$  were calculated to be energetically competitive alternatives to disilene ( $\text{H}_2\text{Si}=\text{SiH}_2$ ).<sup>17d</sup> The spectroscopic identification of  $\text{HSi}(\mu\text{-H})\text{SiH}_2$  (after its theoretical prediction)<sup>17c</sup> and the X-ray structure of a silylene chelate complex<sup>18b</sup> confirm the stability of molecules with divalent silicon. Systematic

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**Figure 2.** Becke3LYP/6-311+G\*\* optimized geometries (distances in angstroms and angles in degrees) for **1–23**.

investigations of silicon–phosphorus compounds (SiH<sub>m</sub>PH<sub>n</sub>,  $m + n = 0–5$ ) at the G2 level also indicate preference for structures with divalent silicon.<sup>19</sup> The multicenter bonded B–H–B and Si–H–Si moieties in boranes and silicon hydrides predicted theoretically<sup>11–13</sup> and observed experimentally<sup>14,18</sup> suggest similar bonding characteristics (B–H–Si) for the hybrid, silaboranes, in addition to structures with low-valent silicon.

Hence, unusual structures and energies of Si<sub>2</sub>B<sub>2</sub>H<sub>4</sub> isomers, compared with boranes, carboranes, and silicon hydrides, can be expected. The parent compounds might be observed in the gas phase or in matrix isolation. Hydrogen-bridged and other structures might be realized in condensed phases when bulky substituents are present. There are numerous precedents.<sup>6,9,14,18</sup> We have employed *ab initio* MO and DFT computations to explore the intriguing structural diversity of the Si<sub>2</sub>B<sub>2</sub>H<sub>4</sub> PES.

### Computational Details

The geometries (Figure 2) were optimized at the electron-correlated MP2(fc)/6-31G\* (MP2)<sup>20a</sup> and Becke3LYP/6-311+G\*\*

(B3LYP)<sup>20c–e</sup> levels using the Gaussian 92/DFT and Gaussian 94 program packages.<sup>21</sup> The nature of the stationary points were characterized by analytically computed frequencies both at MP2 and B3LYP. The total, relative, and zero-point energies are given in Table 1. Relative energies computed at MP2 include scaled (0.95) zero-point vibrational energy corrections.<sup>20b</sup> No scaling was used for the B3LYP zero-point vibrational energies.<sup>22</sup> The stability order of all the structures is comparable at both levels of theory. Hence, only the B3LYP geometries and energies are discussed.<sup>23</sup> Natural population analysis (NPA) charges<sup>24a,b</sup> and Wiberg bond indices (WBIs)<sup>24c,d</sup> for all minimum-energy structures are given in Figure 3.

### Results and Discussion

Nonplanar 1,3-disiladiboretene (C<sub>2v</sub>, **1a**), the most stable Si<sub>2</sub>B<sub>2</sub>H<sub>4</sub> isomer, is puckered<sup>25</sup> by 24.2° and has an inversion barrier of only 1.2 kcal/mol via the D<sub>2h</sub> transition structure **1b**. The analogous 1,3-diboretene (C<sub>2v</sub>) is low in energy, but is not the C<sub>2</sub>B<sub>2</sub>H<sub>4</sub> global minimum.<sup>8</sup> The isovalent and isoelectronic C<sub>4</sub>H<sub>4</sub><sup>2+</sup> and Si<sub>4</sub>H<sub>4</sub><sup>2+</sup> also prefer similar puckered structures

**TABLE 1:** Total (hartrees), Relative (kcal/mol), and Zero-Point Vibrational Energies (ZPE, kcal/mol) for 1–23 at the MP2-fc/6-31G\* (MP2) and BECKE3LYP/6-311+G\*\* (B3LYP) Levels of Optimization (The Number of Imaginary Frequencies (NIM) Are Given in Parentheses)

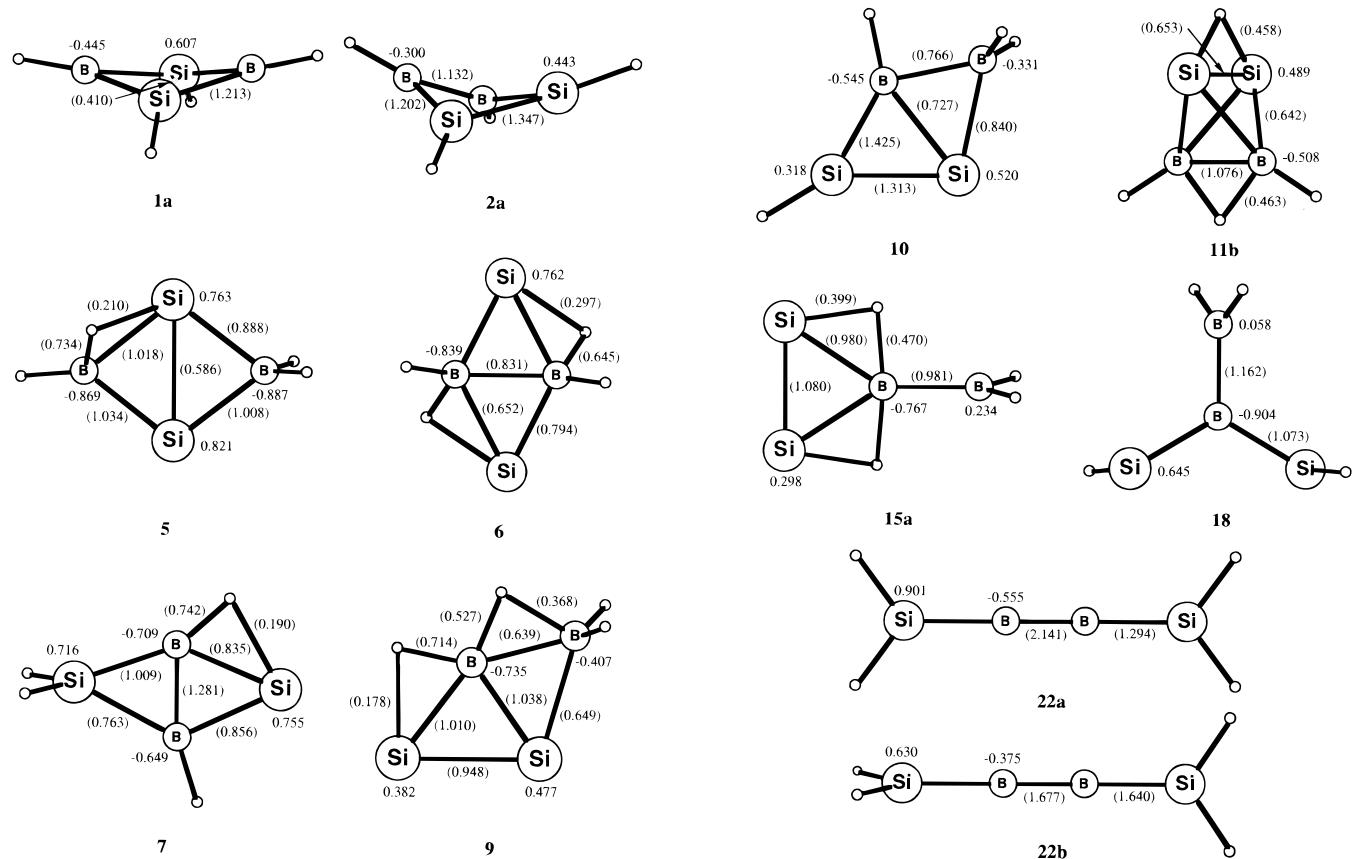
no.	sym.	total energy		relative energy		ZPE (NIM)	
		MP2	B3LYP	MP2	B3LYP	MP2	B3LYP
C <sub>4</sub> H <sub>4</sub> <sup>2+</sup>	D <sub>4h</sub>	-153.343 58	-153.886 52	10.9	8.7	39.0(1)	38.4(1)
C <sub>4</sub> H <sub>4</sub> <sup>2+</sup>	D <sub>2d</sub>	-153.360 80	-153.899 86	0.0	0.0	38.9(0)	38.1(0)
C <sub>2</sub> B <sub>2</sub> H <sub>4</sub>	D <sub>2h</sub>	-127.736 46	-128.264 99	20.7	17.0	35.3(1)	34.5(1)
C <sub>2</sub> B <sub>2</sub> H <sub>4</sub>	C <sub>2v</sub>	-127.769 89	-128.292 46	0.0	0.0	35.6(0)	34.7(0)
Si <sub>4</sub> H <sub>4</sub> <sup>2+</sup>	D <sub>4h</sub>	-1157.528 95	-1159.686 11	-0.2	-0.5	22.4(1)	21.4(1)
Si <sub>4</sub> H <sub>4</sub> <sup>2+</sup>	D <sub>2</sub>	-1157.529 04	-1159.686 70	0.0	0.0	22.7(0)	22.3(0)
<b>1a</b>	C <sub>2v</sub>	-629.747 37	-631.073 66	0.0	0.0	27.0(0)	26.2(0)
<b>1b</b>	D <sub>2h</sub>	-629.743 77	-631.071 38	2.1	1.2	26.8(1)	26.0(1)
<b>2a</b>	C <sub>2</sub>	-629.735 24	-631.063 29	7.3	6.2	26.7(0)	25.9(0)
<b>2b</b>	C <sub>2v</sub>	-629.727 56	-631.059 20	12.0	8.7	26.5(1)	25.8(1)
<b>2c</b>	C <sub>2v</sub>	-629.700 49	-631.026 19	28.3	28.2	25.8(1)	25.2(1)
<b>3a</b>	D <sub>2h</sub>	-629.730 83	-631.059 66	9.3	7.4	25.9(3)	24.8(3)
<b>3b</b>	D <sub>2h</sub>	-629.708 14	-631.035 65	24.6	23.6	27.0(1)	25.9(1)
<b>3c</b>	D <sub>2</sub>	-629.738 11	-631.064 24	5.7	5.3	26.9(1)	25.6(1)
<b>4a</b>	D <sub>2h</sub>	-629.725 64	-631.048 81	13.9	15.4	27.3(1)	26.0(1)
<b>4b</b>	D <sub>2h</sub>	-629.611 74	-630.958 06	82.9	71.0	24.7(3)	24.7(2)
<b>4c</b>	C <sub>2v</sub>	-629.726 00	-631.059 20	13.1	8.2	26.7(1)	25.3(1)
<b>5</b>	C <sub>1</sub>	-629.746 68	-631.070 74	1.3	2.0	27.9(0)	26.4(0)
<b>6</b>	C <sub>2</sub>	-629.735 41	-631.058 10	9.5	10.9	29.1(0)	27.3(0)
<b>7</b>	C <sub>s</sub>	-629.747 04	-631.070 28	0.1	2.0	26.9(0)	26.1(0)
<b>8a</b>	C <sub>2v</sub>	-629.745 02	-631.069 59	0.7	1.9	26.2(1)	25.5(1)
<b>8b</b>	C <sub>2v</sub>	-629.694 06	-631.021 29	32.3	31.9	25.8(2)	25.2(2)
<b>9</b>	C <sub>s</sub>	-629.746 69	-631.064 18	2.0	7.3	28.7(0)	27.6(0)
<b>10</b>	C <sub>1</sub>	-629.734 45	-631.055 08	8.7	11.3	27.6(0)	25.8(0)
<b>11a</b>	C <sub>2v</sub>	-629.646 80	-630.988 23	63.4	53.6	27.3(2)	26.2(2)
<b>11b</b>	C <sub>2v</sub>	-629.729 47	-631.050 61	13.4	16.2	29.3(0)	27.9(0)
<b>11c</b>	C <sub>s</sub>	-629.715 32	-631.044 77	20.8	18.4	27.7(1)	26.5(1)
<b>12</b>	C <sub>2v</sub>	-629.601 34	-630.952 02	90.9	73.4	26.2(3)	23.3(3)
<b>13</b>	C <sub>2v</sub>	-629.657 40	-630.974 50	61.2	61.2	32.0(1)	25.2(1)
<b>14a</b>	C <sub>2v</sub>	-629.686 26	-631.019 66	36.4	32.3	24.9(1)	24.6(1)
<b>14b</b>	C <sub>2v</sub>	-629.615 73	-630.958 33	79.3	70.5	23.5(3)	24.3(2)
<b>15a</b>	C <sub>2v</sub>	-629.701 03	-631.027 84	28.1	27.1	26.0(0)	24.5(0)
<b>15b</b>	C <sub>2v</sub>	-629.697 28	-631.024 62	30.9	29.8	26.4(2)	25.2(2)
<b>16</b>	C <sub>2v</sub>	-629.616 55	-630.951 51	81.2	76.2	26.1(3)	25.7(3)
<b>17</b>	C <sub>2v</sub>	-629.556 81	-630.909 93	115.9	98.8	23.1(4)	22.3(4)
<b>18</b>	C <sub>2</sub>	-629.658 51	-631.009 23	53.8	38.2	24.9(0)	24.0(0)
<b>19a</b>	C <sub>2v</sub>	-629.682 02	-631.019 47	41.1	33.9	27.1(2)	26.1(1)
<b>19b</b>	C <sub>2v</sub>	-629.660 45	-630.998 18	54.2	46.7	26.6(2)	25.5(2)
<b>20</b>	C <sub>2v</sub>	-629.603 41	-630.947 84	87.5	76.2	24.0(3)	23.4(3)
<b>21a</b>	D <sub>2h</sub>	-629.658 77	-630.989 79	55.8	52.2	27.2(2)	25.8(3)
<b>21b</b>	D <sub>2d</sub>	-629.666 90	-631.006 71	50.6	41.8	27.1(2)	26.0(2)
<b>22a</b>	D <sub>2h</sub>	-629.621 57	-630.958 76	79.0	70.5	27.1(0)	24.6(0)
<b>22b</b>	D <sub>2d</sub>	-629.644 69	-630.998 45	62.0	44.9	24.4(0)	23.9(0)
<b>23a</b>	C <sub>2v</sub>	-629.607 57	-630.950 37	85.2	73.8	24.4(2)	22.6(2)
<b>23b</b>	C <sub>2h</sub>	-629.623 69	-630.967 01	77.0	65.6	26.3(1)	24.9(1)

(Figure 1).<sup>1b,7</sup> We have refined the calculations on these isoelectronic and isovalent systems at uniform levels of theory (MP2 and B3LYP) to enable direct comparisons. Table 2 reveals the excellent correlation between the relative energies ( $\Delta E$ ), the puckering angles ( $\Delta\Theta$ ), and the shortening of the transannular distance ( $\Delta D$ ) upon folding. In addition, a comparison of the molecular orbitals (MOs) of the two forms (for example **1a**, **1b** in Scheme 1) shows that the 1,3-antibonding interaction ( $b_{2u}$ ) of the  $\sigma$  orbitals in the planar structure diminishes on puckering ( $b_2$ ).<sup>5a</sup> This stabilizing interaction ( $\Delta\sigma = E_{b2} - E_{b2u}$ ) also follows the trends in relative energies (Table 2). Unlike the  $\sigma$  framework, the  $\pi$  orbital has positive 1,3 overlap ( $b_{3u}$ ), which is destabilized marginally ( $\Delta\pi = E_{a1} - E_{b3u}$ ) upon puckering (Scheme 1). However, the net effect of the  $\sigma$  and  $\pi$  MOs is stabilizing, except for Si<sub>4</sub>H<sub>4</sub><sup>2+</sup>. This qualitative interpretation is reflected in the increased transannular Wiberg bond indices ( $\Delta$ WBI). The NPA charges in **1a** (Si = 0.607; B = -0.445) and in 1,3-diboretene (C = -0.593; B = 0.395) are opposite in sign due to the reversed electronegativity relationship (C = 2.5; B = 2.0; and Si = 1.8).<sup>26</sup>

1,2-Disiladiboretene favors a C<sub>2</sub> minimum **2a** akin to **1a** with a 142.8° B<sup>3</sup>Si<sup>2</sup>B<sup>4</sup>Si<sup>1</sup> dihedral angle. Planar **2b** is the transition

state for ring flipping; the calculated barrier of 2.5 kcal/mol is twice that for the 1,3-disiladiboretene ring inversion (**1a** vs **1b**). The B–B (1.654 Å), B–Si (1.940 Å), and Si–Si (2.196 Å) bonds in **2a** are comparable to the corresponding distances in the model species H<sub>2</sub>B–BH<sub>2</sub> (D<sub>2d</sub>, 1.629 Å), H<sub>2</sub>B–SiH<sub>3</sub> (C<sub>s</sub>, 2.017 Å), and H<sub>2</sub>Si=SiH<sub>2</sub> (C<sub>2h</sub>, 2.173 Å).<sup>27</sup> The WBIs for B–B, B–Si, and Si–Si in **2a** (1.132, 1.202, and 1.347, Figure 2) are greater than 1.0, reflecting the delocalized bonding of the puckered four-membered ring electronic structure. The NPA charges in **2a** (Si = 0.443; B = -0.300), however, are smaller than in **1a**. The 6.2 kcal/mol energy difference between **1a** and **2a** is attributed to the weak  $\pi$ -delocalization in **2a**. Structure **2c**, with short Si–Si and B–B bonds, is a transition state, 28.8 kcal/mol higher in energy than **1a**.

The tendency of silicon to favor divalent structures, e.g. in silicon hydrides and silaboranes,<sup>11–15,17</sup> prompted us to explore similar possibilities on the Si<sub>2</sub>B<sub>2</sub>H<sub>4</sub> PES. Although having short transannular Si–Si and B–B distances, structures **3a**–**3c** and **4a**–**4c** are not minima. The D<sub>2</sub> transition state, **3c**, collapses to the B–H–Si monobridged isomer, **5**, upon relaxation of symmetry constraints. Indeed, **5** is the second most stable Si<sub>2</sub>B<sub>2</sub>H<sub>4</sub> minimum. The B–H (1.280 Å) and H–Si (1.848 Å)



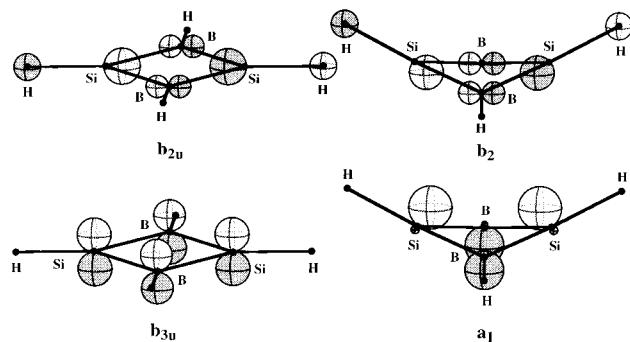
**Figure 3.** Natural population analysis charges and Wiberg bond indices (in parentheses) of Si<sub>2</sub>B<sub>2</sub>H<sub>4</sub> isomers at Becke3LYP/6-311+G\*\* optimized geometries.

**TABLE 2: Changes in Geometries, Energies, and Electronic Structures of 1,3-C<sub>2</sub>B<sub>2</sub>H<sub>4</sub>, C<sub>4</sub>H<sub>4</sub><sup>2+</sup>, 1,3-Si<sub>2</sub>B<sub>2</sub>H<sub>4</sub>, and Si<sub>4</sub>H<sub>4</sub><sup>2+</sup> upon Puckering (at Becke3LYP/6-311+G\*\*\*)**

molecule	$\Delta E^a$	$\Delta\Theta^b$	$\Delta D^c$	$\Delta\sigma^d$	$\Delta\pi^e$	$\Delta\text{WBI}^f$
1,3-C <sub>2</sub> B <sub>2</sub> H <sub>4</sub>	-17.0	51.3	-0.369	-25.1	0.0	0.177
C <sub>4</sub> H <sub>4</sub> <sup>2+</sup>	-8.7	43.7	-0.106	-15.7	0.6	0.037
1,3-Si <sub>2</sub> B <sub>2</sub> H <sub>4</sub>	-1.2	24.2	-0.076	-5.0	3.8	0.016
Si <sub>4</sub> H <sub>4</sub> <sup>2+</sup>	0.5	10.5	-0.005	0.6	3.8	-0.002

<sup>a</sup>  $\Delta E = E_{\text{puckered}} - E_{\text{planar}}$  in kcal/mol. <sup>b</sup>  $\Delta\Theta$  is the puckering angle in degrees (ref 25). <sup>c</sup>  $\Delta D$  is the difference in the 1,3 transannular distance (puckered – planar) in Å. <sup>d</sup>  $\Delta\sigma = [E(b_2) - E(b_{2u})]$  at B3LYP level in kcal/mol (refer to Scheme 1 for symmetry labels). <sup>e</sup>  $\Delta\pi = [E(a_1) - E(b_{3u})]$  at B3LYP level in kcal/mol (refer to Scheme 1 for symmetry labels). <sup>f</sup>  $\Delta\text{WBI} = 1,3\text{-WBI}_{\text{puckered}} - 1,3\text{-WBI}_{\text{planar}}$ .

### SCHEME 1



bridge bond lengths in **5** deviate significantly from the corresponding distances in B<sub>2</sub>H<sub>6</sub> (1.316 Å) and in *cis*-HSi(μ-H)<sub>2</sub>SiH (1.680 Å).<sup>27</sup> The B–H (0.733) and H–Si (0.201) WBIs quantify the B–H–Si bonding interaction (Figure 3). The B–H–Si bridge could be viewed either as a 2c–2e B–H bond hyperconjugating with the p orbital on silicon or, alternatively,

as a 3c–2e bridge polarized toward B–H. Attempts to localize the MOs of **5** using the natural bond orbital (NBO) procedure<sup>24a,b</sup> result in poor descriptions, implying that the bonding is highly delocalized. Although the Si–Si distance in **5** (2.466 Å) is longer than in Si<sub>2</sub>H<sub>6</sub> (2.355 Å),<sup>27</sup> the significant WBI (0.586) shows that **5** is bicyclic. Another bicyclic isomer, **6** (SiBBSi dihedral angle is 131.5°), is 10.9 kcal/mol higher in energy than **1a**. The strong B–B bonding in **6** is evident from the WBI of 0.831, with the B–B distance (1.735 Å) close to the calculated B–B bond length in B<sub>2</sub>H<sub>4</sub> (1.744 Å,  $D_{2h}$ ).<sup>27</sup> A  $\sigma$  lone pair orbital on each silicon and 3c–2e B–H–Si bonding describe part of the electronic structure of **6**. The extent of multicenter bonding is quantified by the WBI data (Figure 3).

The B–H–Si monobridged isomer **7** with planar tetracoordinate boron<sup>28</sup> is related via a B–H wagging motion to the nearly isoenergetic  $C_{2v}$  transition state **8a**.<sup>7b</sup> The all-planar alternative **8b** is a second-order saddle point. The molecular orbitals of **7** include a lone pair on Si and cyclic  $\pi$ -delocalization on the BSiB ring and fail to localize to a Lewis description in terms of NBOs due to the strongly delocalized electronic structure.

As is shown by the significant WBIs, isomer **9** possesses a covalently bound *planar pentacoordinate boron* with unsymmetrical and polarized B–H–Si and B–H–B bridges (Figure 3). The HOMO corresponds to an Si–B–Si  $\pi$  orbital. In addition, a  $\sigma$  lone pair orbital on each silicon and a 3c–2e BBSi bridge is included in the electronic structure description. No Lewis structure is appropriate. When both the bridging hydrogens in **9** become Si–H and B–H terminal bonds, as in **10**, the energy increases by 4.0 kcal/mol. This suggests that the B–H–Si and B–H–B bonds in **9** contribute to the stability of the molecule. The molecular orbitals of **10** feature a lone pair orbital on Si and SiBSi  $\pi$ -delocalization.

The symmetrically hydrogen bridged cyclic structure **11a** possesses two imaginary frequencies, while **11b**, with the Si—H—Si fragment perpendicular to HB—H—BH, is a minimum. Consequently, the B—B and Si—Si distances in **11b** are shortened, but the B—Si bonds are elongated relative to **11a**. The electronic structure of **11b** includes two lone pair MOs and triangulated multicenter bonding patterns. The transition structure **11c** corresponding to the rotation of the —BH<sub>2</sub> moiety is 2.2 kcal/mol above **11b**. Structure **12**, with a symmetrical Si( $\mu$ -H)<sub>2</sub>Si bridge, is a higher order saddle point. In the transition structure **13**, only the bonding combination of the lone pair orbitals is occupied, while the antibonding combination is replaced by the Si—Si  $\pi$  orbital instead.

Interest in the B<sub>2</sub>Si— and Si<sub>2</sub>B— cyclic structures stem from the corresponding carboranes.<sup>8</sup> Krogh-Jesperson *et al.*<sup>29</sup> predicted that 2,3-diboramethylenecyclopropane (2,3-C<sub>2</sub>B<sub>2</sub>H<sub>4</sub>) favors an anti-van't Hoff (perpendicular) over a planar stereochemistry (although neither was found to be a minimum).<sup>8b</sup> Likewise, **14a** and **14b** represent the anti-van't Hoff and van't Hoff stereochemistry among Si<sub>2</sub>B<sub>2</sub>H<sub>4</sub> structures. While **14a** is favored by 38.2 kcal/mol over **14b**, neither is a minimum (Table 1). However, **15a**, with an exocyclic —BH<sub>2</sub> substituent on the Si<sub>2</sub>B ring and a planar pentacoordinate boron, is a minimum. This isomer can be related to the most stable isomer of B<sub>3</sub>H<sub>5</sub><sup>30</sup> considering the isolobal analogy between H—BH<sub>2</sub> and :SiH<sub>2</sub>.<sup>13</sup> The planar conformer, **15b**, is a higher order stationary point. In sharp contrast to the carboranes,<sup>8,29</sup> **15a** prefers bridged B—H—Si bonds over terminal Si—H bonds. The Si—H distance (1.615 Å) is shorter and the B—H (1.393 Å) longer compared to normal bridge bonds,<sup>27</sup> but the MO pattern accounts for the B—H—Si description supported further by the WBIs (Figure 3). The C<sub>2v</sub> structures **16** and **17**, with symmetrical B—H—B and Si—H—Si bridges, are higher order stationary points and hence not discussed further. Thus, all the hydrogen-bridged isomers (Figure 3) exhibit three-center two-electron bonding involving boron and silicon. Such 3c—2e bonding involving three different elements is rare, since the necessary balance in the electronic structure is not often achievable.<sup>31</sup>

Of all the acyclic minima, **18**, with divalent silicones and similar to B<sub>2</sub>H<sub>4</sub> (D<sub>2d</sub>), has the lowest relative energy (38.2 kcal/mol, Table 1). The hyperconjugative interactions between the Si—B bonds and the boron p orbitals are reflected in the short B—B distance (1.615 Å compared to 1.744 Å in D<sub>2h</sub> B<sub>2</sub>H<sub>4</sub>) and the WBI greater than 1.0 (Figure 3). Boryl-substituted disilavinylidene **19a** and **19b** are not minima. Similarly, **20**, analogous to a transition state for the interconversion of 1,2- to 1,3-diboretene on the C<sub>2</sub>B<sub>2</sub>H<sub>4</sub> PES,<sup>8</sup> is a higher order saddle point. The linear structures **21a** and **21b** also are not minima. The corresponding carboranes were predicted to be close in energy to 1,3-diboretene.<sup>8</sup> Substitution of the hydrogen by an —SiH<sub>2</sub> group in B<sub>2</sub>H<sub>2</sub> leads to isomers **22a,b**. The B—B (1.473 Å) and B—Si (1.864 Å) bonds in planar **22a** are comparable to that in HB=BH (1.522 Å) and H<sub>2</sub>Si=BH (1.824 Å).<sup>27</sup> With orthogonal —SiH<sub>2</sub> groups, **22b** is favored by 25.6 kcal/mol over **22a**. The hydrogen-bridged quasi linear structures **23a** and **23b** are high-energy saddle points.

## Conclusions

In contrast to the C<sub>2</sub>B<sub>2</sub>H<sub>4</sub> PES,<sup>8,29</sup> a much larger number of low-energy Si<sub>2</sub>B<sub>2</sub>H<sub>4</sub> minima (**1a**, **2a**, **5–7**, **9**, and **10** are within 12 kcal/mol) has been found. The most stable isomer is 1,3-disiladiboretene (**1a**); the puckered four-membered ring structure is similar to that of 1,3-diboretene.<sup>5,8</sup> However, the computed inversion barrier via the planar 2 $\pi$  aromatic transition state is only 1.2 kcal/mol for **1a**, much lower than the related carborane

barrier (17.0 kcal/mol) at the B3LYP level. B—H—Si bridged isomers **5** and **7**, with divalent silicones, are close in energy to **1a**. Isomer **9**, with a planar pentacoordinate boron, is only 7.3 kcal/mol less stable than **1a**. Isostructural silaborane analogs, **2a**, **8a,b**, **14a,b**, **20**, **21a,b**, and **22a,b**, of the corresponding carboranes are high-energy structures and/or are characterized to be higher order stationary points.

The striking feature of the more stable isomers is their remarkable tendency to possess B—H—Si bridge bonds; C—H—B 3c—2e bonds do not exist among the corresponding carboranes (C<sub>2</sub>B<sub>2</sub>H<sub>4</sub>). The PES of Si<sub>2</sub>B<sub>2</sub>H<sub>4</sub> further demonstrates that isomers with  $\sigma$  lone pairs on silicon and hydrogen bridging (**5–7**, **9**, **11b**, **15a**) are favorable. This contrasts with the preference for multiply bonded structures (cf. **13**, **21b**) exhibited by the analogous carboranes, C<sub>2</sub>B<sub>2</sub>H<sub>4</sub>.<sup>8</sup> Both levels of theory (Becke3LYP/6-311+G\*\* and MP2/6-31G\*) employed here give similar results.

The synthesis of several C<sub>2</sub>B<sub>2</sub>R<sub>4</sub> compounds,<sup>6,9</sup> after their theoretical predictions,<sup>5,29</sup> suggests that experimental efforts toward preparing derivatives of Si<sub>2</sub>B<sub>2</sub>H<sub>4</sub> isomers or gas phase reactions involving silicon hydrides and boranes should be rewarding. Furthermore, examples of H-bridged carborane derivatives<sup>9e</sup> stimulate the search for derivatives of the low-energy H-bridged Si<sub>2</sub>B<sub>2</sub>H<sub>4</sub> isomers (for example **5**, **7**, and **9**), with only the terminal hydrogens replaced by bulky substituents.

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